Kinematics of gravity-capillary waves under an evolving underwater current

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MECHANICAL & AEROSPACE ENGINEERING









Motivation: Remote Sensing and Gravity-Capillary Waves



Farrar et al., 2021. Observations of Ocean Surface Currents.

Studies submesoscale dynamics using these techniques





 $2\lambda sin(\theta) = \lambda_e$



Bodega Marine Laboratory

Some Remote sensing techniques are sensitive to gravity-capillary waves due to Bragg scattering and surface roughness.

Linear Dispersion Relation for Gravity-Capillary waves

Dispersion relation for gravity-capillary waves

$$\omega(k) = \sqrt{gk + \frac{\sigma}{\rho}k^3}$$

$$wavelength$$

$$\lambda = \frac{k}{2\pi}$$

g
ho

How are currents affecting Gravity-Capillary waves?

Let's start considering a **constant velocity** U

Ζ

Х

$$\omega(k) = \sqrt{gk + \frac{\sigma}{\rho}k^3} + \mathbf{u} \cdot \mathbf{k}$$
 Doppler effect

When current flows in same direction as wave, **observed frequency increases**

$$\sigma, \rho$$

In reality, ocean currents vary with depth

 $\omega(k) = \sqrt{gk + \frac{\sigma}{\rho}k^3 + \mathbf{u}\cdot\mathbf{k}}_{\text{Doppler effect}}$ Ζ Х U(z)Which velocity now that U depends on z?

In reality, ocean currents vary with **depth**

$$\omega(k) = \sqrt{gk + \frac{\sigma}{\rho}k^3 + u_{eff}(k)k}$$
Doppler effect
$$U(z)$$
Which velocity now that U
depends on z?

In reality, ocean currents vary with depth

 $\omega(k) = \sqrt{gk + \frac{\sigma}{\rho}k^3 + u_{eff}(k)k} \\ \underset{\text{Doppler effect}}{\text{Doppler effect}}$ Ζ X \mathbf{Z} Waves with different **wave-number** "see" velocities at different depths 7

In reality, ocean currents vary with **depth**

$$\omega(k) = \sqrt{gk + \frac{\sigma}{\rho}k^3} + u_{eff}(k)k$$

Doppler effect
$$\mathbf{U}(\mathbf{z})$$
$$u_{eff}(k) = 2k \int_0^\infty u_L(z)e^{-2kz}dz$$

DNS Two-Phase Flow Setup (Based on Wu et al. 2022)

Waves forced by a turbulent boundary layer



(Wu et al. 2022)

Several non-dimensional parameters

 $u_*/c = 0.25, 0.5, 0.75$ Wind forcing $Re_{\tau} = rac{
ho_a u_* H_a}{\mu_a} = 720$ Turbulent air Reynolds number $k_p H_s = 0.08, 0.16$ Initial wave amplitude $Bo = \frac{(\rho_w - \rho_a)g}{k_n^2 \sigma} = 200$ Wave Bond number $Re_w = \frac{\rho_w c \lambda_p}{\mu_w} = 720$ Wave Reynolds number



Basilisk: Open-source solver (developed by Stephane Popinet): http://basilisk.fr/



Mean flow: Turbulent Boundary Layer in air



Mean flow: Developing viscous layer in water



Developing viscous layer that will transition to a **turbulent boundary layer**

$$U_0(t) = \tau \frac{\Gamma(1)}{\Gamma(3/2)} \frac{\sqrt{\nu_w t}}{\mu_w} = u_*^2 \frac{\rho_a}{\rho_w} \frac{\Gamma(1)}{\Gamma(3/2)} \sqrt{t/\nu_w}$$

(Veron & Melville, 2001)

Mean flow: Turbulent boundary layer in water



 1.Viscous momentum diffusion
 2. Fully developed turbulence

Now we want to look to the space-time spectrum and its evolution





How do energy branches evolve in time?



Branches Time Evolution: Shift from Linear Dispersion Relationship



Branches Time Evolution: Shift from Linear Dispersion Relationship $I = [0, 5] (t/T_p)$ $I = [0, 5] (t/T_p)$ 8--1 $k_f = 1.0k_p$ $\Phi(\omega, k_f, I)/\Phi_{max}$ $k_f = 3.0k_p$ $\Phi(\omega,k,I)/\Phi_{max}$ 6 $k_{f} = 4.0k_{p}$ $a_{m/\alpha}^{d}$ $\sqrt{k \cdot g + \sigma / \rho \cdot k}$ 2N=1-15N=2N=3-13 0. 0 $\hat{2}$ 3 4

6



Branches Time Evolution: First Branch



Branches Time Evolution: Doppler shift first branch



Branches Time Evolution: Doppler shift first branch



Primary mode

$$(k_*, \omega_*)$$

Primary mode

Non-linear (non-resonant) interaction with itself

Higher harmonics

 $(2k_{*}, 2\omega_{*})$

 (k_*, ω_*)

Primary mode

Non-linear (non-resonant) interaction with itself

 (k_*, ω_*)

Higher harmonics

 $(2k_*, 2\omega_*)$ $(3k_*, 3\omega_*)$

Primary mode

Non-linear (non-resonant) interaction with itself

 (k_*, ω_*)

Higher harmonics

 $(2k_*, 2\omega_*)$ $(3k_*, 3\omega_*)$

$$k_N = Nk_*$$
$$\Omega_N = N\omega_*$$

Branches Time Evolution: Higher harmonicsHigher harmonicsPrimary modeNon-linear (non-resonant)
interaction with itselfHigher harmonics
$$(k_*, \omega_*)$$
 $(2k_*, 2\omega_*)$
 $(3k_*, 3\omega_*)$ $(3k_*, 3\omega_*)$
 \downarrow $\Omega_N(k_N) = N\sqrt{gk_N/N + \frac{\sigma}{\rho}(k_N/N)^3}$ $k_N = Nk_*$
 $\Omega_N = N\omega_*$



Branches Time Evolution: NLDR second branch



Branches Time Evolution: NLDR all branches



Branches Time Evolution: phase speed higher modes



Branches Time Evolution: what's the speed of higher modes?



Branches Time Evolution: what's the speed of higher modes?



Branches Time Evolution: higher modes travel same phase speed that primary one



Does the nonlinear dispersion relation hold across different wind and wave initial conditions?

Validation of dispersion relation and depth dependent velocity over different initial conditions



Validation of dispersion relation and depth dependent velocity over different initial conditions



Conclusions

- We performed DNS of fully coupled wind-forced broad-banded wave fields using the open-source solver Basilisk, extending the work of Wu, Popinet, and Deike (2022, JFM).
- The flow is characterized by a turbulent boundary layer on the air side and a viscous self-similar boundary layer in the water, which transitions to a turbulent state.
- Analysis of the evolution of the wavenumber-frequency wave spectra in time:
 - We validate a Non Linear Equation that accounts for the different harmonics:\.
 - **Doppler** shift is characterized with a using a **weighted average depth-varying velocity** integrated from the **DNS**.











Bo = 25

